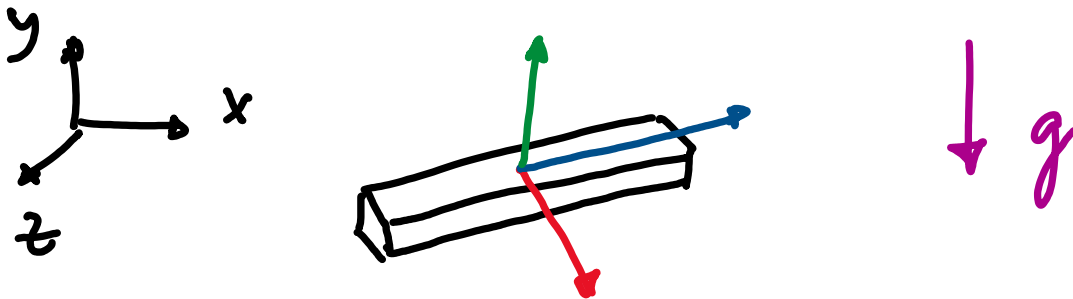


Dynamics

① Free Floating Body (free joint)



Given an initial position & orientation/
initial linear & angular velocity,
describe the movement of the
object

Equations of motion:

(i) Translation :

linear
accel.
eration

$$\left\{ \begin{array}{l} m \dot{v}_x = 0 \\ m \dot{v}_y = 0 \\ m \dot{v}_z = -mg \end{array} \right.$$

3 eqn

$$\left\{ \begin{array}{l} \dot{x} = v_x \\ \dot{y} = v_y \\ \dot{z} = v_z \end{array} \right. \quad \left. \begin{array}{l} \text{velocity} \\ 3 \text{ eqn} \end{array} \right.$$

Written in global frame

(ii) Rotation:

$$\mathbf{I}_b \dot{\boldsymbol{\omega}}_b + \boldsymbol{\omega}_b \times (\mathbf{I}_b \boldsymbol{\omega}_b) = \mathbf{M}_b$$

$\boldsymbol{\omega}_b$ ($\boldsymbol{\Omega}$) - body frame angular velocity

\mathbf{I}_b - body frame inertia

\mathbf{M}_b - Moment in the body frame

$$\dot{\boldsymbol{\omega}}_b = (\mathbf{I}_b)^{-1} [\mathbf{M}_b - \boldsymbol{\omega}_b \times (\mathbf{I}_b \boldsymbol{\omega}_b)]$$

↪ angular acceleration 3 equations

2 methods to compute orientation

(a) Euler angles/rate [1-2-3]

$$\begin{bmatrix} \dot{\phi} \\ \dot{\theta} \\ \dot{\psi} \end{bmatrix} = \frac{1}{\cos \theta} \begin{bmatrix} \cos \psi \cos \theta & \sin \psi \cos \theta & 0 \\ -\sin \psi \cos \theta & \cos \psi \cos \theta & 0 \\ \sin \theta & 0 & 1 \end{bmatrix} \boldsymbol{\omega}_b$$

↪ angular velocity

3 equations

(b) quaternions

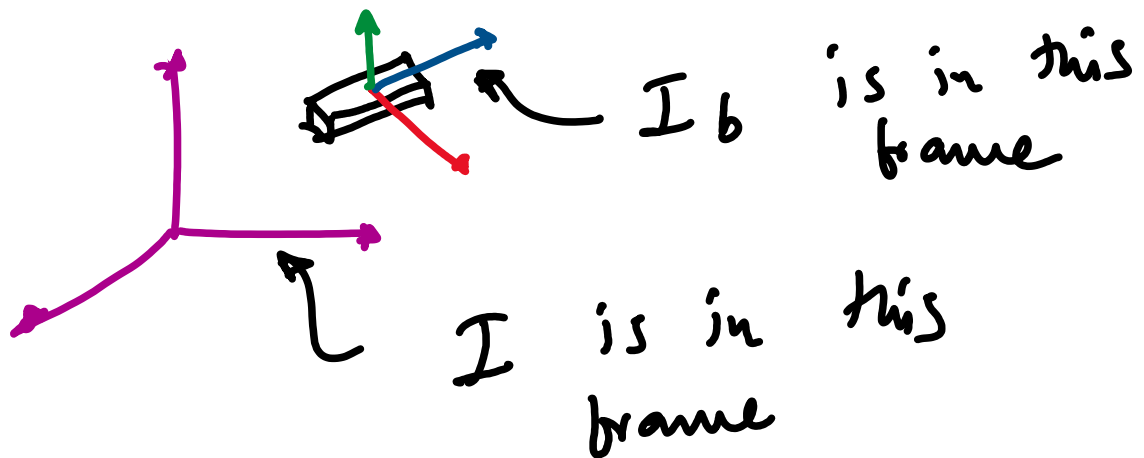
$$\omega_b = 2 \bar{q} \circ \dot{q} \Rightarrow \dot{q} = 0.5 q \circ \omega_b$$

angular velocity

3 equations

I - world frame inertia

I_b - body frame inertia



Relation between I & I_b

$$\begin{aligned}\text{Rotational energy} &= 0.5 \omega^T (I \omega) \\ &= 0.5 \omega_b^T (I_b \omega_b)\end{aligned}$$

$$\omega^T (I \omega) = \underline{\omega_b}^T (I_b \omega_b)$$

But $\omega = R^0 \omega_b = R \omega_b$

$$\omega_b = R^T \omega$$

$$\omega^T (I \omega) = (R^T \omega)^T (I_b R^T \omega)$$

$$\omega^T (I \omega) = \omega^T R (I_b R^T \omega)$$

$$\omega^T (I \omega) = \omega^T R I_b R^T \omega$$

$$I = R I_b R^T$$

"table
wiki moment
of
'inertia'"

Dynamics

2 ways of deriving the equations of motion

① Euler-Lagrange method

- No Free Body Diagram
~ does not give internal forces
- Require Symbolic computation
~ long equations for long chains

② Newton-Euler method

- Requires Free Body Diagram
~ gives interaction forces
 - Symbolic OR Numeric computation
 - symbolic small chains
 - numeric for long chains
- MuJoCo.
- ↪ Recursive Newton Euler Algorithm (RNEA)

① Euler-lagrange method

(1) Write Formula for the position and velocity of the center of mass with respect to the world frame

$$p_c^o = H_i^o p_c^i$$

\uparrow global frame position \uparrow local frame position

$$(i) \quad \underline{v}_c^o = \frac{d}{dt}(p_c^o)$$

$$\underline{\omega}_c^o = \dot{R}_c^o (R_c^o)^T$$

$$(ii) \quad \underline{v}_c^o = J_{v_i} \dot{q}_i$$

$$\underline{\omega}_c^o = \underline{\tilde{J}}_{\omega_i} \dot{q}_i$$

② \mathcal{L} - lagrangian

$$\mathcal{L} = \underline{T} - V$$

$$T = \frac{1}{2} \sum \left[\underbrace{m_i}_{1 \times 3} \underbrace{v_i^T v_i}_{3 \times 1} + \underbrace{\omega_i^T}_{1 \times 3} \underbrace{\left(\underbrace{I_i}_{3 \times 3} \underbrace{\omega_i}_{3 \times 1} \right)}_{3 \times 1} \right]$$

T = kinetic energy

Using the Jacobian we can write

$$T = \frac{1}{2} \dot{q}^T \sum_{i=1}^n \left[m_i (J_{v_i})^T J_{v_i} + J_{w_i}^T \underbrace{(R_i I_b R_i^T)}_{I \text{ (World frame inertia)}} J_{w_i} \right] \dot{q}$$

V - potential energy

$$V = \sum_{i=1}^n m_i g z_G^0 + \frac{1}{2} \sum_{i=1}^m k_{p_i} (r_{p_i} - r_{p_0})^2$$

k_{p_i} - spring constant

r_{p_0}, r_{p_i} - spring length in relaxed state
spring length when stretched

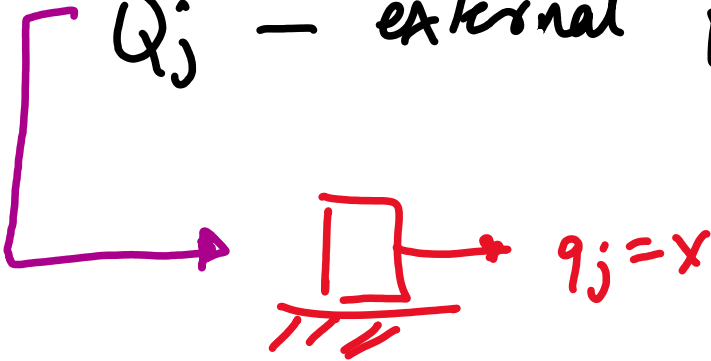
g - gravity (assumed to be along z direction)

③ Equations of motion

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{q}_j} \right) - \frac{\partial \mathcal{L}}{\partial q_j} = Q_j$$

q_j — degree of freedom (e.g. revolute, prismatic)

Q_j — external force / torque (damping friction)



EXAMPLE

Derive the equation of motion for a simple pendulum subject to external torque T_M

①

$$x = l \sin \theta$$

$$y = -l \cos \theta$$

$$\dot{x} = l \cos \theta \dot{\theta}$$

$$\dot{y} = +l \sin \theta \dot{\theta}$$

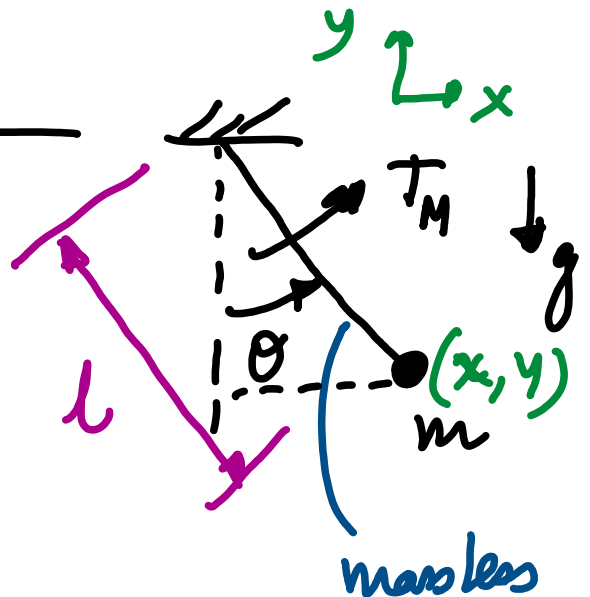
$$\mathbf{v} = \begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} = \begin{pmatrix} l \cos \theta \dot{\theta} \\ l \sin \theta \dot{\theta} \end{pmatrix}$$

② $\mathcal{L} = T - V$

$$= \frac{1}{2} m v^2 - mgy$$

$$= \frac{1}{2} m [l \cos \theta \dot{\theta}, l \sin \theta \dot{\theta}] \begin{pmatrix} l \cos \theta \dot{\theta} \\ l \sin \theta \dot{\theta} \end{pmatrix} + mgl \cos \theta$$

$$\mathcal{L} = \frac{1}{2} m l^2 \dot{\theta}^2 + mgl \cos \theta$$



$$\mathcal{L} = \frac{1}{2} m l^2 \ddot{\theta}^2 + m g l \cos \theta$$

$$\textcircled{3} \quad \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\theta}} \right) - \frac{\partial \mathcal{L}}{\partial \theta} = T_m$$

$$\frac{d}{dt} \left[\frac{1}{2} m l^2 (2\dot{\theta}) \right] - m g l (-\sin \theta) = T_m$$

$$m l^2 \ddot{\theta} + m g l \sin \theta = T_m$$

check: $T_m = 0$

$$m l^2 \ddot{\theta} + m g l \sin \theta = 0$$

$$\ddot{\theta} = -\frac{g}{l} \sin \theta$$

eqⁿ of
a simple
pendulum

General form of Equation of motion
for a manipulator

$$\rightarrow \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{q}_i} \right) - \frac{\partial \mathcal{L}}{\partial q_i} = Q_i$$

↓ Spong Chapter 6

torque
↓

$$\rightarrow \boxed{\underbrace{M(q)}_{n \times n} \underbrace{\ddot{q}}_{n \times 1} + \underbrace{C(q, \dot{q})}_{n \times 1} + \underbrace{G(q, g)}_{n \times 1} = U}$$

q

$M(q)$ mass matrix,
depends only on q

$C(q, \dot{q})$ non-linear coriolis / centripetal
acceleration term.
depends on q, \dot{q}

$G(q, g)$ gravity term, it depends
on gravity, g

$$x = \underline{l \sin \theta}$$

$$\dot{x} = l \cos \theta \dot{\theta}$$

another way

$$\begin{aligned} \dot{x} &= \left(\frac{\partial x}{\partial \theta} \right) \dot{\theta} \quad \rightarrow \text{chain rule} \\ &= (l \cos \theta) \dot{\theta} \quad \text{jacobian} \end{aligned}$$

$$EOM = \begin{Bmatrix} EOM_0 \\ EOM_1 \\ EOM_2 \end{Bmatrix} \quad \left\{ \begin{array}{l} M, C, G = ? \end{array} \right.$$

$$(1) \quad M = \frac{\partial EOM}{\partial \ddot{\theta}} \quad \ddot{\theta} = [\ddot{\theta}_1, \ddot{\theta}_2, \ddot{\theta}_3]$$

$$(2) \quad C = -EOM (\dot{q} = 0, \dot{q}, q, q = 0)$$

$$(3) \quad G = -EOM (\ddot{q} = 0, \dot{q} = 0, q, q)$$

$$M \ddot{q} + C + G = \vec{0}$$

$$\ddot{q} = -M^{-1}(C+G)$$

Simulate : integrati~

$$\checkmark \quad \underline{\dot{q}} = - \int \underline{M^{-1}(C+G)} dt = \underline{\omega}$$

$$\underline{q} = \int \underline{\omega} dt$$

angular
accelerations

once we obtain q we can
do an animation $\setminus [q_1, q_2, q_3]$

odeint (runge-kutta adaptive)

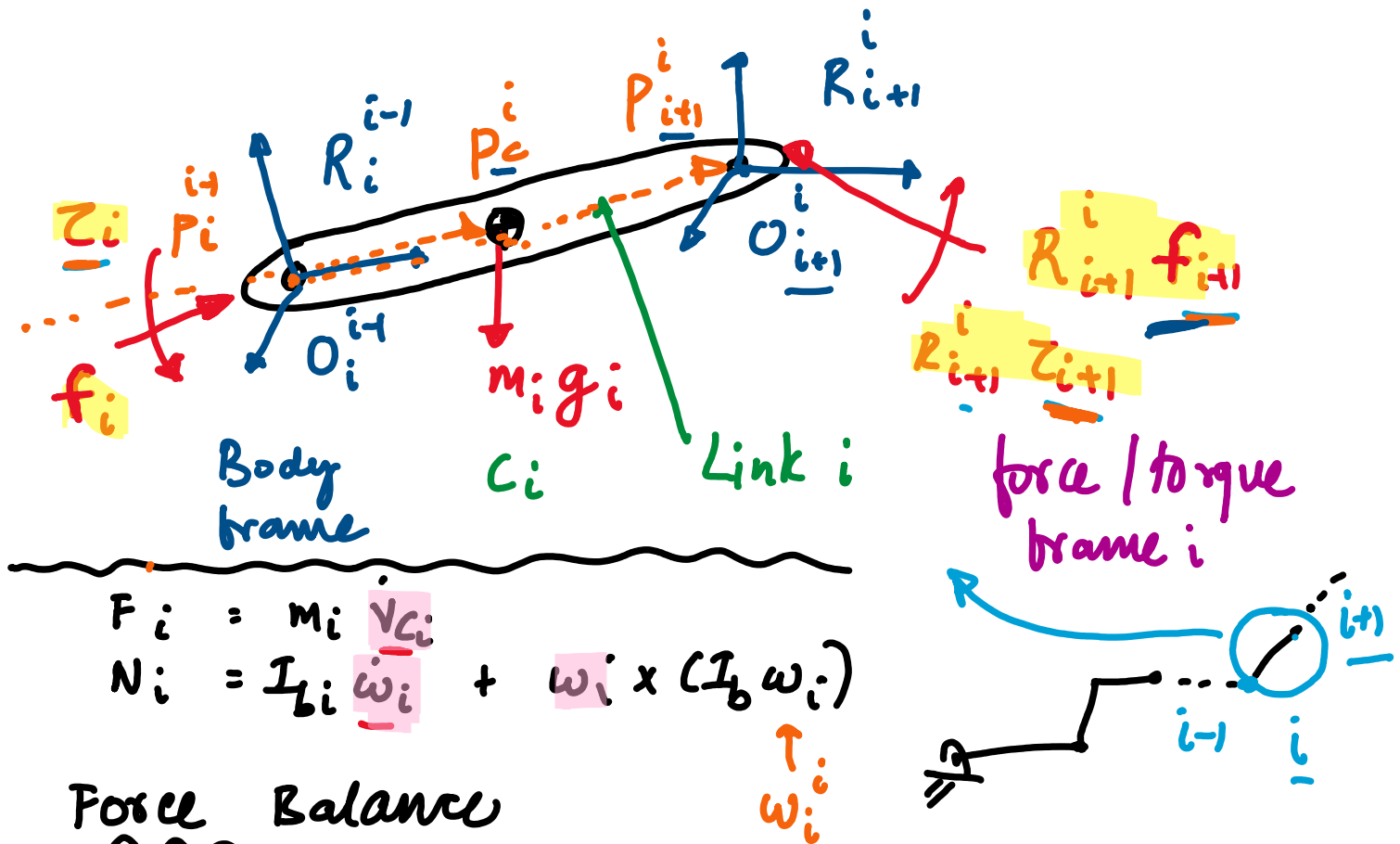
① Euler's method

② Runge kutta fixed, order 4

MuJoCo

more accurate

II Newton - Euler Method (Recursive Newton-Euler Algorithm)



$$F_i = m_i \dot{v}_{C_i}$$

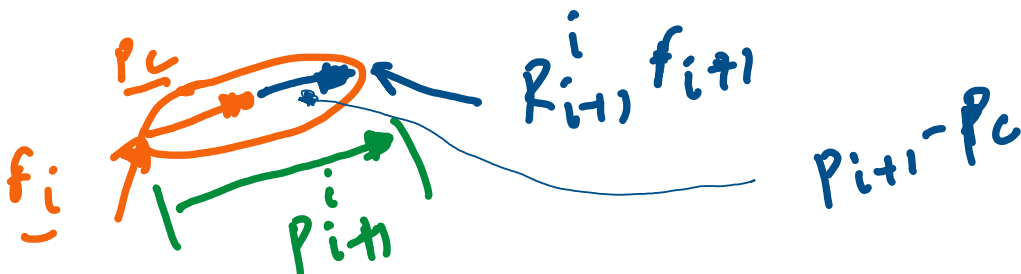
$$N_i = I_{b_i} \dot{\omega}_i + \omega_i \times (I_b \omega_i)$$

Force Balance

$$f_i = F_i + R_{i+1}^i f_{i+1} - m_i g_i$$

Torque Balance about point C_i

$$\checkmark \quad \tau_i = N_i + R_{i+1}^i \tau_{i+1} + P_c^i \times f_i + (P_{i+1}^i - P_c^i) \times R_{i+1}^i f_{i+1}$$



Velocities and Accelerations

Joint i is **revolute**

$$\underline{\omega}_i = (R_i^{i-1})^T \underline{\omega}_{i-1} + \dot{q}_i \hat{n}_i \quad \leftarrow \text{same as velocity.}$$

$$\underline{\dot{\omega}}_i = (R_i^{i-1})^T \underline{\dot{\omega}}_{i-1} + \omega_i \times \dot{q}_i \hat{n}_i + \ddot{q}_i \hat{n}_i$$

$$\underline{\dot{v}}_i = (R_i^{i-1})^T \left[\underline{\dot{v}}_{i-1} + \underline{\dot{\omega}}_{i-1} \times p_i^{i-1} + \omega_{i-1} \times (\omega_{i-1} \times p_i^{i-1}) \right]$$

\ddot{O}_i^{i-1}

$$\underline{\dot{v}}_{c_i} = \underline{\dot{v}}_i + \underline{\dot{\omega}}_i \times p_c^i + \omega_i \times (\omega_i \times p_c^i)$$

Joint i is **prismatic**

$$v = v_0 + \omega \times (v_0 \times r) + \dot{\omega} \times r$$

$$\omega_i = (R_i^{i-1})^T \omega_{i-1}$$

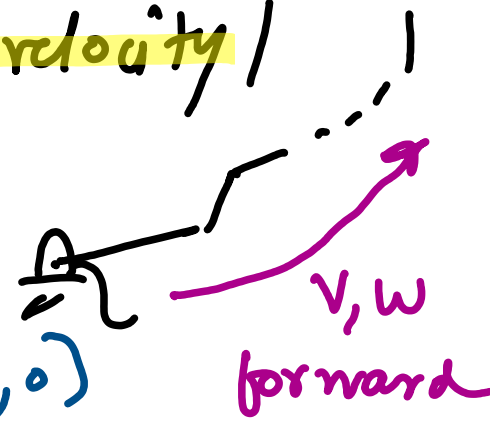
$$\dot{\omega}_i = (R_i^{i-1})^T \dot{\omega}_{i-1}$$

$$\underline{\dot{v}}_i = (R_i^{i-1})^T \left[\underline{\dot{v}}_{i-1} + \underline{\dot{\omega}}_{i-1} \times p_i^{i-1} + \omega_{i-1} \times (\omega_{i-1} \times p_i^{i-1}) + \ddot{O}_i^{i-1} + \dot{q}_i \hat{n}_{i-1} + 2\omega_{i-1} \times \dot{q}_{i-1} \hat{n}_{i-1} \right]$$

$$\underline{\dot{v}}_{c_i} = \underline{\dot{v}}_i + \underline{\dot{\omega}}_i \times p_c^{i-1} + \omega_{i-1} \times (\omega_{i-1} \times p_c^{i-1}) + \ddot{q}_i \hat{n}_i + 2\omega_i \times \dot{q}_i \hat{n}_i$$

How to use these equations?

① Forward recursion for velocity / acceleration

$$V_0 = [0, 0, 0] \quad \dot{V}_0 = [0, 0, 0] \quad \omega_0 = [0, 0, 0] \quad \dot{\omega}_0 = [0, 0, 0]$$


Use recursion to compute

$$V_i, \dot{V}_i, \omega_i, \dot{\omega}_i, \dots$$

$$V_i, \omega_i, \dot{V}_i, \dot{\omega}_i = f(V_{i-1}, \omega_{i-1}, \dot{V}_{i-1}, \dot{\omega}_{i-1})$$

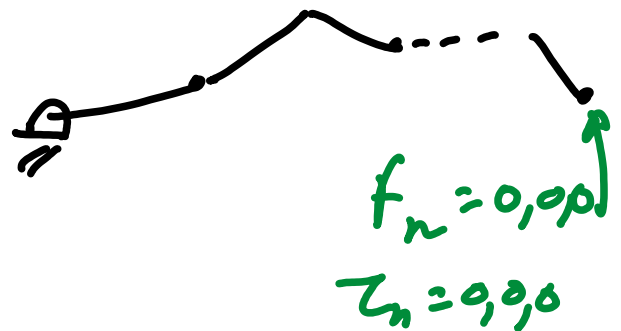
$i=1$

② Backward recursion for forces / torques

$$F_n = [0, 0, 0]$$

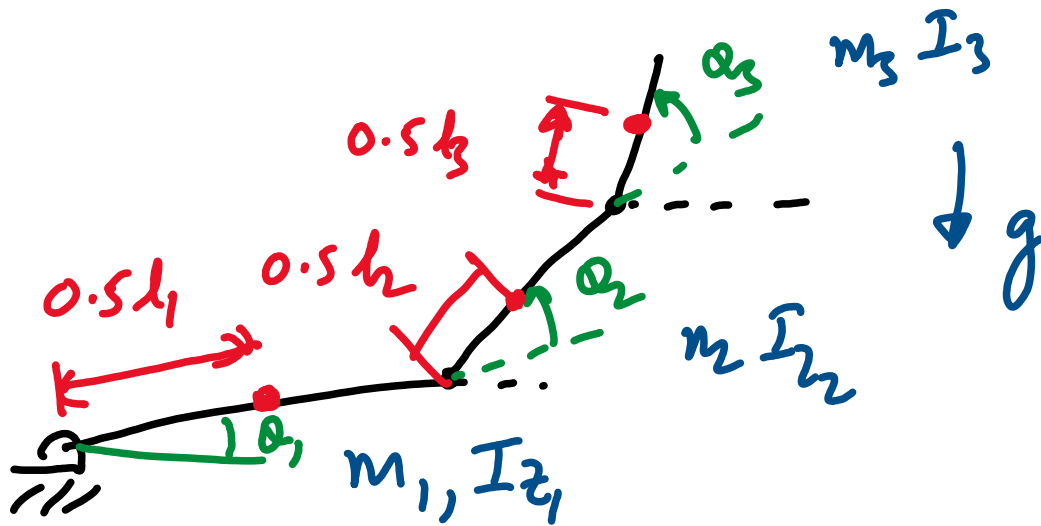
$$Z_n = [0, 0, 0]$$

$$\underline{F_i}, \underline{Z_i} = \bar{F}(F_{i+1}, Z_{i+1})$$



$$\text{COM: } Z_1(\hat{n}_1), Z_2(\hat{n}_2), \dots \quad \hat{n}_i = \text{axis of motion for } q_i$$

Derive the equations of motion of a 3-link planar manipulator



Compute M numerically

$$\textcircled{M} \ddot{q} + C(q, \dot{q}) + G(q, g) = \underline{z}$$

numerically

$$\begin{bmatrix} M_{11} & M_{12} & M_{13} \\ M_{21} & M_{22} & M_{23} \\ M_{31} & M_{32} & M_{33} \end{bmatrix} \begin{bmatrix} \ddot{q}_1 \\ \ddot{q}_2 \\ \ddot{q}_3 \end{bmatrix} + C(q, \dot{q}) + G(q, g) = \underline{z}$$

$$\rightarrow \dot{q} = 0, g = 0 \Rightarrow C = 0, G = 0$$

$$\text{Set } \ddot{q}_1 = 1 ; \ddot{q}_2 = \ddot{q}_3 = 0$$

$$\underline{z} = \begin{bmatrix} M_{11} \\ M_{21} \\ M_{31} \end{bmatrix}$$

$$\text{Set } \ddot{q}_2 = 1 ; \ddot{q}_1 = \ddot{q}_3 = 0$$

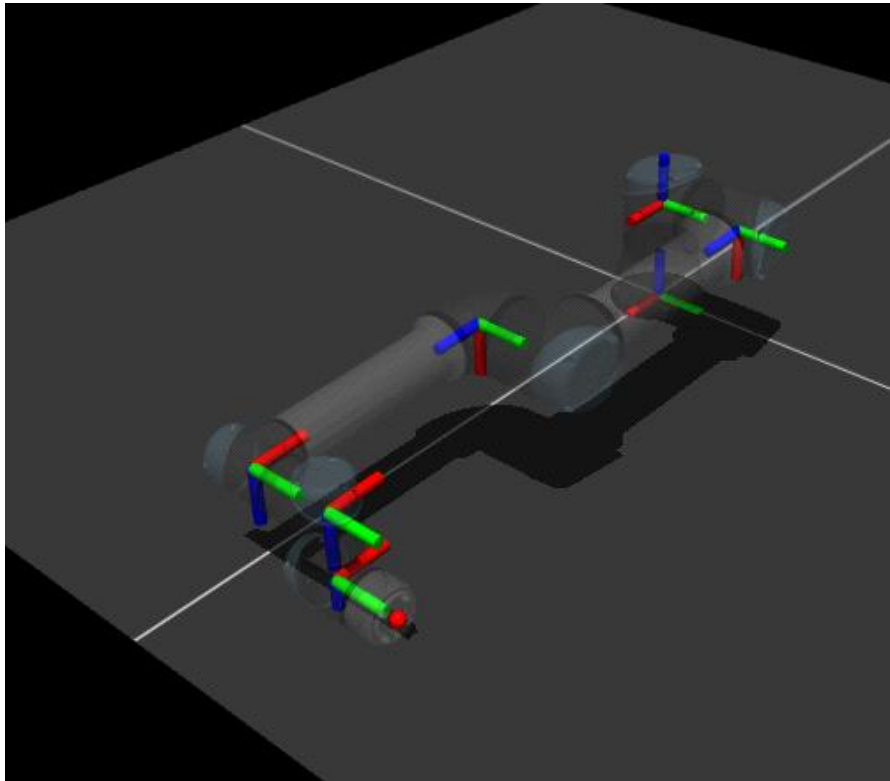
$$\underline{z} = \begin{bmatrix} M_{12} \\ M_{22} \\ M_{32} \end{bmatrix}$$

$$\text{Set } \ddot{q}_3 = 1 ; \ddot{q}_1 = \ddot{q}_2 = 0$$

$$\underline{z} = \begin{bmatrix} M_{13} \\ M_{23} \\ M_{33} \end{bmatrix}$$

3 calls
to
forward
backward
reduction

Recursive Newton Euler Algorithm for UR5



$$M(q)\ddot{q} + C(q, \dot{q}) + G(q, g) = \tau$$

$$q \in \mathbb{R}^6 \quad C, G \in \mathbb{R}^6 \quad M \in \mathbb{R}^{6 \times 6}$$

- ① Euler-Lagrange \sim symbolic (lengthy)
- ② Newton-Euler \sim " (")
- ③ Recursive Newton-Euler Algorithm (RNEA) \sim numerical

Use TMT method to compute M

$$M = T^T \bar{M} T$$

$$M = J^T \bar{M} J$$

6×6 6×36 36×36 36×6

J - jacobian of the center of mass of each link concatenated such that $J \in \mathbb{R}^{36 \times 6}$

$$\bar{M} = \begin{bmatrix} M_1 & 0 & 0 & 0 & 0 & 0 \\ 0 & \bar{I}_1 & 0 & 0 & 0 & 0 \\ 0 & 0 & M_2 & 0 & 0 & 0 \\ 0 & 0 & 0 & \bar{I}_2 & 0 & 0 \\ & & & & \ddots & \\ & & & & & M_6 \\ & & & & & 0 & \bar{I}_6 \end{bmatrix}$$

36×36

$$M_i = \begin{bmatrix} m_i & 0 & 0 \\ 0 & m_i & 0 \\ 0 & 0 & m_i \end{bmatrix}$$

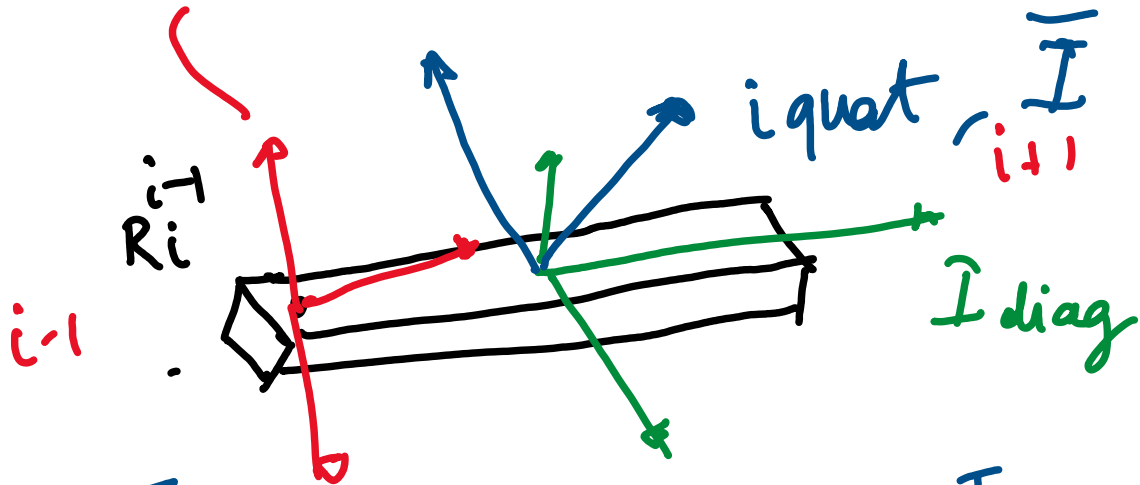
↑

$$\bar{I}_i = R_i^{i-1} \text{Rinertia} I_{diag} R_{inertia}^T R_i^{i-1}$$

$$\text{Rinertia} = \text{iquat2 Rotation}$$

↑

$\langle \text{inertial} \quad \text{quat} = \quad \text{diag} = \begin{bmatrix} - & - & - \end{bmatrix} \rangle$
 body frame
 iquat
 I diag



$$\bar{I} = R_{\text{inertia}} I_{\text{diag}} R_{\text{inertia}}^T$$

$$I_i = R_i^{i-1} \bar{I} (R_i^{i-1})^T$$

$R_{\text{inertia}} = \text{ran. quat 2 rotation (iquat)}$